

Critical couplings and string tensions from two-lattice matching of RG decimations

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Critical couplings and string tensions in $SU(2)$ and $SU(3)$ lattice gauge theory are calculated by two-lattice matching of RG block transformations. The transformations are of the potential moving type generating plaquette actions with large number of group characters and exhibit rapid approach to a unique renormalized trajectory. Fixing the critical coupling $\beta(N_\tau)$ at one value of temporal lattice extent N_τ by MC simulation, the critical couplings for other values of N_τ are then obtained by lattice matching of these decimations. $\beta(N_\tau)$ values are thus obtained for a range of N_τ and found to be in agreement with MC simulation results to within a few percent in all cases. A similar procedure allows the calculation of string tensions for a range of β values with similarly good agreement with MC data.

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1. Introduction

The renormalization group (RG) based technique of “lattice matching” via block transformations relates physical quantities on different lattices. It thus provides a method for computation of a physical quantity at different lattice spacings (couplings). To apply it one needs to implement RG block transformations on the approach to the Wilsonian Renormalized Trajectory (RT). This can be done in various ways. One way is numerical implementation of the RG blocking by Monte Carlo RG (MCRG) techniques. This is the method that has mostly been used in the literature. Another approach is to implement blocking by explicit RG recursion relations that can, to varying degree, be explicitly carried out by numerical-analytical means. This is the method followed here. Specifically, we employ explicit RG recursion relations of the “potential moving” type. These block transformations (decimations) are of course approximate but can, in principle, be systematically improved. They turn out to be surprisingly effective. In the following we apply lattice matching of these decimations to obtain critical couplings and string tensions for the $SU(2)$ and $SU(3)$ pure gauge theories. A more detailed account has appeared in [1].

2. RG blocking recursions and lattice matching

RG blocking recursions We start by assuming a general plaquette action $A_p(U_p, n)$ on lattice of spacing $b^n a$ given in terms of the character expansion of its exponential:

$$\exp(-A_p(U_p, n)) = \sum_j d_j F_j(n) \chi_j(U_p). \quad (2.1)$$

The sum is over all inequivalent irreducible representations labeled by j , with characters χ_j of dimension d_j . The action itself is, of course, completely specified by the set of $F_j(n)$ coefficients, and vice versa, and of the general form:

$$A_p(U_p, n) = \sum_j \frac{1}{d_j} \beta_j(n) \frac{1}{2l_j} [\chi_j(U_p) + \chi_j(U_p^{-1})] \quad (2.2)$$

with $l_j = 1$ for self-conjugate and $l_j = 2$ for non-self-conjugate representations. (For $SU(2)$, in particular, $l_j = 1$ for all j .)

It is useful to define an effective coupling $g^{(n)}$ characterizing a given action of the form (2.2):

$$\beta^{(n)} = \frac{2N}{g^{(n)2}} \equiv 2N \frac{d^2 A_p(e^{i\theta\hat{m}\cdot t}, n)}{d\theta^2} \Big|_{\theta=0}. \quad (2.3)$$

Here $\{t\}$ are the $SU(N)$ generators and \hat{m} a unit vector. ((2.3) is of course independent of the direction \hat{m}). In the perturbative regime this reduces to the usual definition of gauge coupling. In the non-perturbative regime any definition of a ‘coupling’ is of course some scheme-dependent choice. We adopt (2.3) to track the RG recursion flows; it provides an efficient parametrization of the renormalized trajectory below.

The lattice block step $b^n a \rightarrow b^{n+1} a$ may now be formulated as a prescription for the character expansion coefficients $F_j(n+1)$ in terms of the $F_j(n)$ ’s:

$$F_j(n+1) = \left[\int dU \left[\sum_k d_k F_k(n) \chi_k(U) \right]^{\zeta^{(d-2)}} \frac{1}{d_j} \chi_j^*(U) \right]^{r^2}. \quad (2.4)$$

To complete the prescription we must specify the renormalization parameters ζ, r . We take

$$\zeta = b \left[1 - c g^{(n)2} \right] \quad (2.5)$$

$$r = b \left[1 - c g^{(n)2} \right] \quad (2.6)$$

with c an adjustable decimation parameter to be tuned for optimization as explained below. It is convenient to work with normalized coefficients $f_j = F_j/F_0$ by factoring out the trivial representation coefficient in (2.1). Effective couplings (2.3) are also conveniently computed directly in terms of the $\{f_j\}$.

RG flows and lattice matching Consider now a general lattice system described by an action $A(K)$ with set of couplings $K = \{K_i\}$. Successive RG blockings by a scale factor b generate a flow in action space: $K \rightarrow K^{(1)} \rightarrow K^{(2)} \rightarrow \dots \rightarrow K^{(n)} \rightarrow \dots$, where $K^{(n)} = \{K_i^{(n)}\}$ denotes the couplings after n blocking steps. If flows from K and K' reach the same point on RT after n and n' steps, then

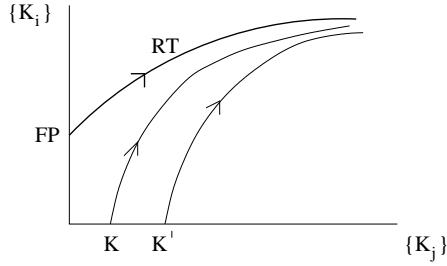


Figure 1: Flows towards the RT from two different starting couplings K and K' .

the corresponding lattice correlation lengths ξ, ξ' and spacings a, a' are related as

$$\xi' = b^{-(n-n')} \xi \quad a' = b^{(n-n')} a. \quad (2.7)$$

To identify such pairs of couplings we need ascertain that after n and n' RG steps, respectively, the same point is reached on the RT. This can be done in either of two ways: (i) show that the corresponding actions coincide: $A(K^{(n)}) = A(K'^{(n')})$. This requires that one obtain the blocked action at each step; or (ii) show that the expectations of every operator, measured after performing the corresponding number of blocking steps from the initial two actions, agree. Either way, identifying such pairs (K, n) , and (K', n') is referred to as two-lattice matching [2].

If blockings are performed numerically by MCRG, the second method appears easier to use. Obtaining the blocked action can be difficult, whereas it is possible, at least in principle, to generate a Boltzmann-weighted configuration ensemble for the blocked action by instead blocking the configurations of an ensemble generated from the original action. These can then be used to measure observables [3]. In practice, of course, due to lattice size limitations, only a rather small number of block steps is possible by MCRG, so getting close enough to the RT is not guaranteed. As a general observation, the location of the fixed point being block definition dependent, appropriate fine-tuning of any decimation free parameters can be crucial for achieving rapid approach in few steps.

Here we employ two-lattice matching with RG block transformations implemented by the recursions (2.4) described above. They can be explicitly evaluated to any desired accuracy on lattices of any size, so no inherent limitations due to finite size arise. The blocked action resulting after each RG step is explicitly obtained, so it can be used to ascertain approach to the RT and perform two-lattice matching. The transformations contain one parameter (cf. (2.5) - (2.6)), which should be fixed for matching optimization.

A basic feature of our decimations is that, regardless of the choice of the initial plaquette action, a single step suffices to generate an action of the form (2.2) generally containing a large (infinite) set of representations. This is important as flow in such a large-dimensional interaction space makes it possible to avoid getting stuck at (finite-dimensional) lattice artifact boundaries. Furthermore, MCRG construction of blocked actions [4] shows that one-plaquette terms with a large number of characters are the most relevant action terms for long-scale dynamics. This is precisely the type of action resulting from our decimations, and may be the reason for their apparent efficacy in computing long-distance dynamics observables as seen below.

In the following the starting action for our decimations ($n = 0$) will always be taken to be the fundamental representation Wilson action. (Other choices such as mixed actions containing several representations can be treated in exactly the same way.) One finds that the flow under successive decimations reaches a unique RT irrespective of such a choice, though of course the number of steps needed to reach it depends on the initial point in action space. With the fundamental Wilson action as the starting action the approach is found to be very rapid as illustrated in Fig. 2.

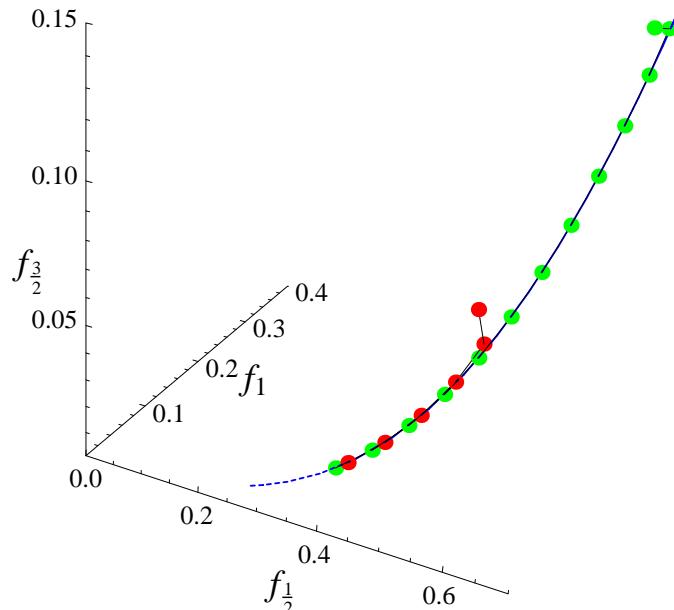


Figure 2: Flow from the $SU(2)$ fund. Wilson action with $\beta = 4$ (green dots) and $\beta = 2.5$ (red dots) showing rapid approach to the RT. First three non-trivial (normalized) expansion coefficients shown.

The effective coupling (2.3) provides a good way to label points along the RT. If, starting from

some Wilson action coupling β , after n steps the point $\beta^{(n)}(\beta)$ lies on the RT, subsequent RG steps generate a sequence of points $\beta^{(n+1)}, \beta^{(n+2)}, \dots$ hopping along the RT. With scale factor $b = 2$, and for all large and intermediate values of $\beta^{(n)}$, the effective beta function is varying slowly enough for a linear interpolation to provide an excellent approximation to the RT points lying between pairs of neighboring points $\beta^{(n)}, \beta^{(n+1)}$.

The decimations become exact in the $\beta \rightarrow \infty$ limit. Computation of the step scaling function (beta function) from the decimations in the weak coupling scaling region reproduces the perturbation theory prediction to within 2% – 3%. We next use them to obtain critical couplings and string tensions by lattice matching for the $SU(2)$ and $SU(3)$ gauge theories.

3. Critical couplings and string couplings by two-lattice matching

3.1 Critical couplings

At physical temperature $T = 1/aN_\tau$, lattice with temporal extent N_τ , spacing a and lattice with N'_τ , a' are related by:

$$a' = \frac{N_\tau}{N'_\tau} a. \quad (3.1)$$

If after blocking n and n' times, respectively, the two flows reach the same point on the RT, using (2.7) this implies

$$n - n' = \log_b \left(\frac{N_\tau}{N'_\tau} \right). \quad (3.2)$$

So, at $T = T_c$ one has

$$\beta^{(n)}(\beta_c(N_\tau)) = \beta^{(n')}(\beta_c(N'_\tau)). \quad (3.3)$$

This suggests the following simple matching procedure.

(i) Assuming $\beta_c(N_\tau)$ known for one N_τ , take

$$\begin{aligned} n &= \log_b N_\tau + m \\ n' &= \log_b N'_\tau + m \end{aligned} \quad (3.4)$$

Integer $m = 0, 1, \dots$ is chosen so n, n' large enough to be on the RT. (If the so-chosen n or/and n' turn out to be non-integer, one performs $[n]$ and $[n] + 1$ steps, where $[n]$ is the nearest integer to the chosen n from below, and uses interpolation for the RT points in-between as mentioned above.)

(ii) With n, n' and $\beta_c(N_\tau)$ given, solve (3.3) for $\beta_c(N'_\tau)$. This means that the starting point of the flow on the N'_τ lattice is adjusted to satisfy (3.3).

3.2 String tensions ($T = 0$)

A similar procedure allows one to obtain string tensions by matching. Assume that two RG flows from starting Wilson action at β_0 and β_1 end up at the same RT point after n_0 and n_1 steps, respectively. Then one has

$$\beta^{(n_0)}(\beta_0) = \beta^{(n_1)}(\beta_1) \quad (3.5)$$

and

$$a_1 \sqrt{\sigma} = b^{(n_0 - n_1)} a_0 \sqrt{\sigma} \quad (3.6)$$

Suppose $a_0 \sqrt{\sigma}$ known. Choose n_0 large enough to be on the RT. Determine n_1 so that (3.5) is satisfied. $a_1 \sqrt{\sigma}$ is then obtained directly from (3.6).

3.3 Numerical results

It is important to maintain high accuracy in working with the expansions (2.1) under blocking iteration. For $SU(2)$ we typically use fifty group characters in the expansions (2.1). This implies for, say, $\beta = 5$ omitted higher character coefficients $f_j = F_j/F_0$, and accompanying bounds on the series remainder, of the order of 10^{-45} . For $SU(3)$ we truncate (2.1) at characters $j \equiv (p, q)$ with $p \geq 20, q \geq 20$; this implies remainders at $\beta = 10$ of less than 10^{-12} .

The scale factor is always taken to be $b = 2$. The adjustable parameter in the decimation recursions (2.4) - (2.6) is c , which we set at $c = 0.10$ in the case of $SU(2)$ and $c = 0.24$ in the case of $SU(3)$. With no other parameters present, straightforward numerical evaluation of the recursion relations can then be carried out.

We take one value of $\beta_c(N_\tau)$ from MC data, which serves to fix the scale and apply the procedure above to obtain critical coupling values for other lattices. Results for $SU(2)$ are shown in Table 1. Two sets of computed β_c values are shown in Table 1 (columns 1 and 2) corresponding to two different choices of the MC data point (underlined entries). The table also shows comparison with the values obtained by MC simulation [5] - [7] in each case (column 3). The agreement is very good - at the 2% - 3% level. Results for critical couplings in the $SU(3)$ gauge theory are displayed in Table 2. Agreement with MC simulation data [5] is again very good, typically within a few percent.

N_τ	β_c	β_c	$\beta_c(\text{MC})$
3	2.1875	2.1957	2.1768(30)
4	2.2909	<u>2.2991</u>	2.2991(02)
5	2.3600	2.3683	2.3726(45)
6	2.4175	2.4258	2.4265(30)
8	2.5097	2.5180	2.5104(02)
12	<u>2.6355</u>	2.6440	2.6355(10)
16	2.7275	2.7361	2.7310(20)
32	2.9487	2.9574	

Table 1: Critical couplings $\beta_c(N_\tau)$ for $SU(2)$ computed from lattice matching of decimations. Column 1 and 2 show the values obtained for two different choices (underlined entries) of the one data point taken from MC data (see text). Column 3 shows the values from MC simulations for comparison.

String tensions in $SU(3)$ obtained by the method above are displayed in the same format in Table 3. Very similar results are obtained for $SU(2)$ [1]. Good agreement with MC data [7] - [9] is again obtained in all cases.

Consideration of fermionic observables by similar RG recursion methods is a rather more demanding proposition. Some preliminary attempts are reported in [10].

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N_τ	β_c	β_c	$\beta_c(\text{MC})$
4	5.6501	5.6329	5.6925(002)
6	<u>5.8941</u>	5.8773	5.8941(005)
8	6.0773	6.0595	6.0010(250), 6.0625(18)
10	6.2018	6.1837	6.1600(70)
12	6.3084	6.2900	6.2680(120), 6.3385(55)
14	6.4015	<u>6.3830</u>	6.3830(100)
16	6.4845	6.4658	6.4500(500)
32	6.9024	6.8829	

Table 2: Critical couplings $\beta_c(N_\tau)$ for $SU(3)$ computed from lattice matching of decimations and comparison with MC simulation data. Same format as in Table 1.

β	$a\sqrt{\sigma}$	$a\sqrt{\sigma}$	$a\sqrt{\sigma}(\text{MC})$
5.54	0.5580	0.5878	0.5727(52)
5.6	0.5070	<u>0.5295</u>	0.5295(09), 0.5064(28)
5.7	0.4205	0.4264	0.4099(12), 0.3879(39)
5.8	0.3486	0.3508	0.3302(15)
5.9	0.2919	0.2931	0.2702(19)
6.0	0.2465	0.2433	0.2269(62), 0.2209(23)
6.2	0.1698	0.1671	0.1619(19), 0.1604(11)
6.4	<u>0.1214</u>	0.1180	0.1214(12), 0.1218(28)
6.5	0.1010	0.0983	0.1068(09)
6.8	0.0616	0.0599	0.0738(20)

Table 3: String tensions $a\sqrt{\sigma}$ for $SU(3)$ computed from lattice matching of decimations. Same format as in Table 1.

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